EARLY STRUCTURE FORMATION FROM PRIMORDIAL DENSITY FLUCTUATIONS WITH A BLUE-TILTED POWER SPECTRUM

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ABSTRACT

While observations of large-scale structure and the cosmic microwave background (CMB) provide a strong constraint on the amplitude of the primordial power spectrum (PPS) on scales larger than 10 Mpc, the amplitude of the power spectrum on sub-galactic length scales is much more poorly constrained. We study early structure formation in a cosmological model with a blue-tilted PPS. We assume that the standard scale-invariant PPS is modified at small length scales as \( P(k) \sim k^{m_s} \) with \( m_s > 1 \). We run a series of cosmological hydrodynamic simulations to examine the dependence of the formation epoch and the characteristic mass of primordial stars on the tilt of the PPS. In models with \( m_s > 1 \), star-forming gas clouds are formed at \( z > 100 \), when formation of hydrogen molecules is inefficient because of the intense CMB radiation. Without efficient coolant, the gas clouds gravitationally contract while keeping a high temperature. The protostars formed in such “hot” clouds grow very rapidly by accretion to become extremely massive stars that may leave massive black holes with a few hundred solar-masses at \( z > 100 \). The shape of the PPS critically affects the properties and the formation epoch of the first generation of stars. Future experiments of the CMB polarization and the spectrum distortion may provide important information on the nature of the first stars and their formation epoch, and hence on the shape of the small-scale power spectrum.

Subject headings: cosmology: theory – dark ages, reionization, primordial stars – stars: Population III – stars: formation – method: numerical

1. INTRODUCTION

The formation of the first stars marks the end of the cosmic Dark Ages, when their energetic photons initiate reionization of the inter-galactic medium. The exact epoch when cosmic reionization began is still uncertain, but evidence is mounting from an array of observations that the first luminous objects such as galaxies and quasars were formed very early (Watson et al. 2013; Wu et al. 2015; Planck Collaboration XIII 2013; Planck Collaboration XIII 2015). Unfortunately, such tight constraints are derived only at large length scales, up to the wavenumber of \( k \sim 0.2 \text{ Mpc}^{-1} \) (Hlozek et al. 2012). Most of theoretical studies thus assume a scale-invariant PPS and adopt significant extrapolation of the observationally determined large-scale power spectrum. The observational evidence on sub-galactic scales is confusing: milli-lensing measurements suggest substructure consistent with the extrapolation of the Λ-Cold Dark Matter (Λ-CDM) power spectrum to smaller scales (Hezaveh et al. 2013; Bussmann et al. 2013; Xu et al. 2013) yet studies of local dwarf galaxies suggest significant problems with Λ-CDM on small scales (Strigari et al. 2007). It is not only interesting to explore how early structure formation is affected by enhanced density fluctuations at small scales, but also timely because the existence of supermassive black holes at very high redshifts suggests rapid growth of small scale structure in the early universe.

A variety of possibilities are proposed from the physics of the early universe that posit scale-dependent PPS. Popular single-field inflation models predict a nearly scale-free and slightly red power spectrum for adiabatic scalar perturbations. There are also variant models that predict scale-dependent power...
spectrum with kinks, bumps, and other types of features (e.g., Starobinski 1992, Adams et al. 1997, Kawaguchi et al. 2008, Biswas et al. 2010, Barnaby 2010). Of particular interest to us in the present paper are the models that generate density perturbations with a blue-tilded power spectrum or with enhancement at small length scales (Covi & Lyth 1999, Martin & Brandenberger 2001, Gong & Sasaki 2011). It is naively expected that structure forms early in such models, yielding abundant small mass dark halos. Because primordial star formation involves a number of physical processes such as chemistry and radiative transfer as well as gravitational assembly of dark halos, the initial density perturbations or the PPS can affect the formation of primordial stars in a complicated manner. Clearly, it is important to study the properties of the first structure using cosmological simulations with the relevant physics included.

In the present paper, we perform a series of cosmological simulations for different PPS models to study early structure formation in detail. In particular, we examine how the slope of the PPS at the small length scales affects the formation epoch and the mass of the first generation of stars. We begin by describing the calculation methods in Section 2. Section 3 shows the simulation results of primordial star formation for different PPS models. We discuss the dependence of the primordial star formation on the PPS model in Section 4. Throughout this paper, we adopt the standard Λ-CDM cosmology with the total matter density $\Omega_m = 0.3086$, the baryon density $\Omega_b = 0.04825$, the dark energy density $\Omega_{\Lambda} = 0.6914$ in units of the critical density, the Hubble constant $h = 0.6777$, and the primordial index $n_s = 0.9611$ (Planck Collaboration XVI 2014).

2. NUMERICAL SIMULATIONS

We perform a series of cosmological simulations that start from the initial conditions generated for different PPS models $P_{\text{prim}}(k)$. In the next sections, we describe the numerical methods with emphasis on the improvement over our earlier numerical simulations (Hirano et al. 2014, 2015).

2.1. Initial Conditions

The standard scale-invariant PPS is given by

$$P_{\text{prim}}(k) \propto k^{n_s},$$  \hspace{1cm} (1)

whereas the model power spectrum with enhancement at small scales is given by

$$P_{\text{prim}}(k) \propto k^{n_s} \left\{ \begin{array}{ll} 1 & (k < k_p) \\ k^{s-m_s} & (k > k_p) \end{array} \right.$$  \hspace{1cm} (2)

We adopt two parameters that characterize the spectrum; $k_p$ is the pivot wavenumber above which the spectral index is set to be $n_s$. We consider $k_p = 100, 300, \text{ and } 500 \ h \ Mpc^{-1}$ and $m_s = 1.5, 2, \text{ and } 2.5$.

Figure 1 shows the matter power spectra for models with $k_p = 100 \ h \ Mpc^{-1}$. The fiducial pure power-law PPS case is also shown there for comparison. We use the transfer function of Eisenstein & Hu (1999) for this plot and also for generating the initial conditions using the publicly available code MUSIC (Hahn & Abel 2011). The density fluctuation amplitude is normalized by $\sigma_8$. Because our simulation volume is very small, we set $\sigma_8$ larger than the current standard value of about 0.8 (Planck Collaboration XIII 2015), so that we can simulate the formation of high-$\sigma$ peak halos. In practice, we set $\sigma_8 = 1.5, 2, \text{ and } 2.5$. In total, we generate 18 cosmological initial conditions with the parameter sets given in Table 1.

The effect of the enhanced small-scale power is clearly seen in the mass function (Fig. 2). The halo mass functions differ by orders of magnitude at $z = 100$. Figure 2 also shows that the Press-Schechter mass function agrees fairly well with the simulation results (see also Reed et al. 2005).
2.2. Cosmological Simulations

We perform the cosmological simulations using the parallel N-body / Smoothed Particle Hydrodynamics (SPH) code Gadget-2 (Springel 2005) suitably modified for the primordial star formation simulations (Yoshida et al. 2002, 2007; Hirano et al. 2014, 2015). In order to achieve sufficient numerical resolution, we use hierarchical zoom initial conditions. The parent simulation volume is \( L_{\text{box}} = 100 \ h^{-1} \ \text{kpc} \) on a side. The mass of the refined dark matter particles in the zoomed region is \( 0.80 \ M_\odot \), and that of the baryonic component is \( 0.15 \ M_\odot \). We follow structure formation from \( z_{\text{ini}} = 499 \) until the central hydrogen number density reaches \( n_{H,\text{cen}} \approx 10^7 \ \text{cm}^{-3} \) in a gas cloud that first collapses (see Section 2.3). The initial ionization fraction is computed by using RECFAST (Seager et al. 1999, 2000; Wong et al. 2008) as \( x_e = 6.88 \times 10^{-4} \) at the initial redshift. During the cloud collapse, we keep the refinement criterion that the local Jeans length is always resolved by 15 times the local smoothing length by progressively increasing the spatial resolution by the particle-splitting technique (Kitsionas & Whitworth 2002).

Let us briefly describe the overall trend found from the simulation result. In the standard model, the first stars are formed at \( z \sim 20 - 50 \) (e.g., Tegmark et al. 1997; Reed et al. 2005; Yoshida et al. 2003). In one of our simulations with \( \sigma_8 = 1.5 \), the first star is formed at \( z_{\text{form}} = 24 \) within the simulation volume of 100 \( h^{-1} \ \text{kpc} \) on a side. We label this sample as ID 100. Here, the first digit “1” indicates the value of \( \sigma_8 \) (see Table 1). We also find that, by changing \( \sigma_8 \) to 2.0 and 2.5 (ID = 200 and 300), the first star in the simulation volume is formed earlier at \( z_{\text{form}} = 35 \) and 46, respectively.

2.2.1. \( H_2 \) Formation in the Early Universe

We are interested in the thermal evolution of a primordial gas cloud formed at \( z > 100 \). Because CMB photons at such high redshifts are capable of dissociating molecules and ions, several chemical reactions need to be included in order to follow the thermal and chemical evolution accurately.

At \( z < 100 \), the main \( H_2 \) formation path is the so-called \( H^- \) channel:

\[
H + e^- \rightarrow H^- + \nu, \quad (3)
\]

\[
H^- + H \rightarrow H_2 + e^- \quad (4)
\]

Because \( H^- \) ions are destroyed by CMB photons at \( z > 100 \) via the photo-detachment reaction

\[
H^- + \nu \rightarrow H + e^- \quad (5)
\]

the \( H^- \) channel is not effective at such early epochs. An alternative path is the \( H_2^+ \) channel:

\[
H^+ + H \rightarrow H_2^+ + \nu \quad (6)
\]

\[
H_2^+ + H \rightarrow H_3^+ + H^+ \quad (7)
\]

An opposing dissociative process by energetic photons operates as

\[
H_2^+ + \nu \rightarrow H^+ + H \quad (8)
\]

At \( z > 120 \), most of the hydrogen molecules are formed by the \( H_2^+ \) channel. The mean \( H_2 \) fraction in the Universe reaches \( 10^{-6} \) at \( z \sim 60 \). We update the chemical network by introducing the above reactions (see a review by Galli & Palla 2013). The reaction rates are based on Coppola et al. (2013). We describe more details of the chemistry implementation in Appendix A.

2.3. Final Stellar Masses

It is important to study the characteristic mass of the first stars in our variant cosmological models with blue-tilted PPS. We calculate the mass of a primordial star by using two different methods. One is direct calculation of gas accretion onto the central protostar us-
Fig. 3.—Projected densities for dark matter (Panel a) and baryonic components (b). We also plot the projected gas temperature (c) at the formation sites of the primordial stars ID 300 (top; $z_{\text{form}} = 46$) and ID 313 (bottom; $z_{\text{form}} = 186$). The plotted region has a proper 100 pc on a side. We scale the density in units of particle number density per cubic-centimeter.

ing the two-dimensional axisymmetric radiation hydrodynamic (RHD) simulation coupled with the stellar evolution (Hosokawa et al. 2011, 2012; Hirano et al. 2014, 2015). We perform the calculations for 6 cases. We label them as ID = 100, 113, 200, 213, 300, and 313. Here, the second digit of the object ID indicates the pivot wavenumber, and the last digit indicates the value of ms (see Table 1). For each gas cloud, we generate the initial conditions for the RHD simulation by suitably mapping the final snapshot of our cosmological simulation at the time when the density of the collapsing cloud core reaches $n_{\text{H, cen}} = 10^{13}$ cm$^{-3}$. The RHD simulations are run until the mass accretion rate onto the central star falls below $10^{-4}$ $M_\odot$ yr$^{-1}$. At this point the stellar growth is almost completely halted and the resultant mass can be regarded as the final stellar mass. We denote the stellar masses calculated directly by the RHD simulation as $M_{\ast, \text{cal}}$.

We also estimate the final stellar masses by using a fitting function that is derived from a number of RHD simulations of accreting protostars in our earlier studies. The fitting function essentially correlates the stellar mass $M_{\ast, \text{est}}$ to the gas infall rate $M_{\text{Jeans}}$ evaluated for a gravitationally Jeans unstable gas cloud. It is given by Hirano et al. (2015) as

$$M_{\ast, \text{est}} = 250 \left( \frac{M_{\text{Jeans}}}{2.8 \times 10^{-3} M_\odot \text{yr}^{-1}} \right)^{0.7} M_\odot. \quad (9)$$

By adopting this formula, we can skip the numerically costly RHD simulation in the accretion phase. Note that the difference between the calculated and estimated values is about fifty percent, and mostly within a factor of two (see fig. 14b in Hirano et al. 2014). Although this might still appear insufficient to determine the stellar mass for all the cases in Table 1, the mass variations caused by model parameters ($k_p$, $m_s$, etc) are typically larger than the uncertainty in the estimated mass (see Table 1). Because we aim at investigating how cosmological parameters such as the slope of the power spectrum affect the formation of early halos and gas clouds, it would be clearer to interpret the resulting stellar masses in terms of gas clouds' properties such as the Jeans mass and accretion rate. Therefore, we use the estimated stellar masses (Eq. 9) for all the 18 cases to discuss the qualitative dependence on the adopted PPS models.

Our stellar mass estimates are based on the results of radiation-hydrodynamics simulations which follow only the evolution of the central star. Recent three-dimensional simulations show that the circumstellar disk can gravitationally fragment to produce multiple clumps (sink particles; e.g., Clark et al. 2011; Greif et al. 2011, 2012; Stacy & Bromm 2013). Such fragments, if some of them survive over a long period of disk accretion, would form multiple stars with correspondingly small masses. High-resolution simulations show also that fragments efficiently merge onto the central star in roughly an orbital timescale (e.g., Greif et al. 2012; Vorobyov et al. 2013). In the latter case, the final stellar mass would be similar with the calculated value by our radiation-hydrodynamics simulation. Conservatively, we interpret the derived $M_\ast$ as the total mass of stars formed, rather than the mass of a single star.

3. RESULTS

3.1. Major Effects
We find major effects of the small-scale power spectrum on the properties of host halos and on the star formation within them. Figure 3 shows the local conditions around the formation site of the primordial star. The top panels are for the case with normal scale-invariant PPS (ID 300). The primordial star-forming cloud is formed at \( z_{\text{form}} = 46 \) inside a dark matter mini-halo that has a virial radius of \( R_{\text{vir}} \sim 19 \) pc. The bottom panels are for the case with an extremely enhanced PPS (ID 313). The gas cloud is formed earlier at \( z_{\text{form}} = 186 \) inside a more compact mini-halo with \( R_{\text{vir}} \sim 3.6 \) pc. Figure 3(a) shows the density distribution of dark matter component within a proper 100 pc on a side. Note that the mean density is higher in the bottom case than in the top because of the cosmic expansion, as \( \rho \propto (1 + z)^{3} \). The number of small mass dark matter clumps is significantly larger in the bottom panel, reflecting the enhanced initial density perturbations at the corresponding length scales (e.g., Figs. 1 and 2). Interestingly, we find a secondary star-forming cloud formed in the vicinity of the first one within several 10 pc (Fig. 3(b)). Such a close pair of primordial stars likely affect each other through radiative, dynamical, and chemical feedback effects. In the top panels, we also find a neighboring mini-halo but it has not grown sufficiently to host a dense star-forming cloud (see also the temperature distribution in Panel c).

The host mini-halos forming at very high redshifts are generally more compact and have larger dark matter densities, whereas the characteristic density of the gas cloud is determined largely by molecular hydrogen cooling. Consequently, the baryonic component becomes dynamically dominant only after the gas density exceeds \( \sim 10^{5} \) cm\(^{-3}\) in very early halos.

### 3.2. Formation Epoch and \( \text{H}_2 \) Formation

Figure 4(a) shows the density evolution of the gas cloud’s core as a function of redshift for the cases with \( \sigma_{8} = 2.5 \). With more enhanced power (larger \( m_s \)), the clouds collapse earlier. In the extreme case with \( m_s = 2.5 \) (red line), the cloud collapses at \( z \sim 186 \). Then \( \text{H}_2 \) formation mainly proceeds via the \( \text{H}_2^{+} \) channel, but the photo-dissociation by CMB prevents \( \text{H}_2 \) formation. Hence the molecular fraction remains low for ID 313 until three-body \( \text{H}_2 \) formation operates at very high densities (Fig. 4(b)). The formation of \( \text{H}_2 \) molecules critically affects the thermal evolution, prestellar collapse, and the subsequent accretion process onto the central protostar.

### 3.3. Thermal and Dynamical Properties of Clouds

Figure 5 shows the evolution and the properties of collapsing clouds for the six cases for which direct RHD simulations are performed. We use black lines for models with the standard power-law spectrum \( (m_s = n_s) \) and we use red lines for models with enhanced small scale power \( (m_s > 1) \). The thermal evolution in the standard model is well understood and can be described as follows (see Panel a). The temperature first increases due to gravitational contraction but then decreases owing to \( \text{H}_2 \) cooling, when the molecular fraction reaches the critical value that is needed to cool the gas within a Hubble time \( f_{\text{H}_2} \sim 5 \times 10^{-4} \) (see Fig. 5), e.g., Tegmark et al. (1997). Note that the minimum temperature is floored by \( T_{\text{CMB}} \) at high redshifts. The gas temperature starts increasing again when the compressional heating overcomes the radiative cooling, leading to the formation of a protostar.

In our models with enhanced small scale power (red lines), the cloud evolves on a higher temperature track because of the lower \( \text{H}_2 \) fraction (Panel b). With larger \( m_s \), the density at which the \( \text{H}_2 \) fraction reaches the critical value for cooling is larger. Interestingly, for ID 313 (red solid line), the \( \text{H}_2 \) formation fraction reaches the critical value until three-body \( \text{H}_2 \) formation occurs at \( n_{\text{H}_2, \text{cen}} > 10^{8} \) cm\(^{-3}\). In this case, the temperature at the onset of run-away collapse is 700 K, much higher than that found in the usual primordial star formation. One can naively expect that protostars formed in such a ‘hot’ cloud will grow very rapidly.

The stellar mass can be estimated from Eq. (10). Effectively, the mass accretion history determines the protostellar evolution and the strength of radiative feedback that halts the accretion of surrounding medium. Figure 5(c) shows the gas infall rate as a function of the
enclosed gas mass. On average, the gas infall rate systematically increases at higher redshift (see the difference from long-dashed to solid lines). However, the gas infall rate does not monotonically increase with \( m_\text{vir} \). This is because the host mini-halos of the gas clouds are formed at earlier epochs but contain less amount of gas.

Figure 5(d) shows the ratio of rotational speed to the Keplerian speed. Interestingly the gas clouds formed at earlier epochs spin significantly more slowly in the models with enhanced power. This may reflect the effective redistribution of the (orbital) angular momenta brought in by numerous minor mergers. Stars formed in a gas cloud with less spin tend to be more massive because of relatively rapid spherical accretion onto the protostar (Hirano et al. 2014).

3.4. Dependence on Model Parameters

It is important to study the effect of the pivot scale \( k_\text{p} \) of the PPS on gas cloud formation. There is an important physical scale that characterizes primordial star formation. Let us denote the characteristic wavenumber \( k_\text{h} \) for a mini-halo that has a physical size \( R_\text{h} \) so that \( k_\text{h} \sim 2\pi/R_\text{h} \). When \( k_\text{p} \) is smaller than \( k_\text{h} \), the enhanced power generates more structure of mini-halo sizes.

In the fiducial case of ID = 100, the host halo has a virial mass of \( \sim 2 \times 10^5 M_\odot \). The corresponding wavenumber is \( k_\text{h} = 2\pi/R_\text{h} \sim 500 \text{ h Mpc}^{-1} \). When \( k_\text{p} \) is comparable to \( k_\text{h} \), as in the cases with \( k_\text{p} = 300 \) and \( 500 \text{ h Mpc}^{-1} \), increasing \( m_\text{s} \) does not significantly affect \( z_\text{form} \) and \( M_\text{vir} \) (see Table 1). This is in direct contrast to the trend found when \( k_\text{p} < k_\text{h} \); increasing \( m_\text{s} \) actually resulted in larger \( z_\text{form} \) and smaller \( M_\text{vir} \). This can be understood as follows. The formation epoch is largely determined by the amplitude of large-scale density perturbations whereas the enhanced power promotes formation of small mass halos. Therefore, when \( k_\text{p} \) is small, density perturbations at all the relevant length scales are enhanced. The mini-halos are then formed early, and are physically compact.

It is worth considering possible halo-to-halo variation. Our understanding so far is based on a limited number of samples. Hirano et al. (2014, 2015) show substantial differences in the properties of star-forming clouds even with similar masses and similar formation epochs that are in turn determined largely by the cosmological parameters. In principle, using a much larger number of halos could show the general trend and variations more clearly, but we have focused on qualitative differences caused by the shape of PPS.

4. DISCUSSION

We have studied how the shape of the primordial power spectrum at small scales affects early structure formation. For a significantly enhanced PPS, star-forming gas clouds are formed at extremely early epochs of \( z > 100 \) (ID 213 and 313), when the usual gas-phase reaction of molecular hydrogen formation does not operate. At \( 100 < z < 300 \), hydrogen molecules are formed via the \( H_2 \) channel and thus the gas can still cool and condense to form primordial stars. The formed stars have masses as large as \( \sim 300 M_\odot \).

At even earlier epochs of \( z > 300 \), there is no efficient formation channel of \( H_2 \) and the possible path to
star formation is through cooling by atomic hydrogen. If the highly dense and compact clumps are formed in a similar manner to the so-called direct collapse black hole formation model (e.g., Inayoshi et al. 2014). Such very massive stars gravitationally collapse to form remnant massive black hole at the end of their lives, leaving promising “seeds” for the formation of the recently discovered supermassive black hole (SMBH; Wu et al. 2014).

There is an additional generic physical process that can significantly affect the structure formation in the early universe; the relative motions between dark matter and baryons caused by acoustic oscillations at the cosmic recombination era (Tseliakhovich & Hirata 2010). The relative streaming motions prevent the gas contraction into dark matter mini-halo and delay the star formation, and thus possibly change the physical condition of the first star formation (see Fialkov 2014 for recent review). The effect is important at the very high redshifts considered in this paper, since the root-mean-square streaming velocity scales with redshift as \( \propto (1 + z) \). We have run several additional simulations and confirmed that, under significant streaming motions, primordial gas clouds are formed later in large host halos than presented in the present paper. There is an interesting possibility that streaming motions delay the first star formation to the extent that gas collapse is suppressed in mini-halos and more massive halos eventually host very massive stars, which can become the seed of SMBH (Tanaka & Li 2014). The blue-tilted primordial power-spectral can effectively increase the number of massive halos in the early universe. Further results from simulations with the streaming motions will be presented elsewhere (Hirano et al., in preparation).

Observationally, an independent probe of the PPS at small length scales might be provided by the existence of small mass dark halos in and around galaxies. Following the analytical model of Zhao et al. (2003), we estimate that an extremely large number of halos with masses of \( \sim 1 M_\odot \) are formed for \( m_s = 1.5 \) and \( k_p = 100 \text{ Mpc}^{-1} \). The typical formation epoch of such solar-mass halos is \( z \sim 300 \), and thus the halos may remain as very compact objects. However, even the dense and compact clumps can be destroyed by tidal stripping/disruption and encounter with stars in the Galaxy. It is thus difficult to estimate the abundance of the mini- and micro-halos surviving until today, either by large scale simulation (e.g., Diemand et al. 2005; Green & Goodwin 2007) or by analytical calculation (Zhao et al. 2005). Nevertheless, if signatures of small mass dark halos are detected through, for instance, dark matter annihilation or direct detection experiments, one can obtain invaluable information on their origin, i.e., the small-scale power spectrum.

Currently, there are little direct probes of the small-scale primordial density fluctuations, and only model dependent constraints are available (Chluba et al. 2012; Berezinsky et al. 2014; Natarajan et al. 2015). We look forward to future CMB experiments such as PRISM and PIXIE (Kogut et al. 2011) that can infer the formation epoch and the spectral energy distribution (hence roughly the typical mass) of the first stars through measurement of CMB spectral distortions and the detection of low level of early reionization at \( l > 10 \) (e.g., Ricotti et al. 2003). Finally, James Webb Space Telescope (Windhorst et al. 2009) and WFIRST observations (Spiegel et al. 2013) will probe the existence of quasars at even higher redshifts. Understanding the nature of the first stars may thus provide a route to probe observationally the small-scale primordial power spectrum.

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A. REACTION RATE CALCULATION OF H$_2^+$ PHOTODISSOCIATION

The reaction rate calculation of H$_2^+$ photodissociation is highly uncertain. In our previous calculations [Hirano et al. 2014, 2015], we adopt the fitting function of Galli & Palla (1998) assuming the LTE-level population. The accurate reaction rate can be calculated by considering the H$_2^+$-level population (Hirata & Padmanabhan 2006; Coppola et al. 2011). The time evolution of H$_2$ abundance at $z > 60$, which is the relevant epoch in our study, differs significantly depending on the adopted reaction rate (see fig. 3a in Galli & Palla 2013). We take a straightforward approach by comparing simulations with different reaction rates. Fortunately, we find that the uncertainty in the reaction rate gives essentially no influence on the primordial star formation. Stars are formed in dense gas clouds in which a critical amount of hydrogen molecules are formed to overcome the compressional heating by their radiative cooling. When a gas cloud becomes gravitationally unstable after the so-called loitering phase, the H$_2$ fraction within the cloud reaches a threshold value that is independent on the treatment of H$_2^+$ photodissociation. We thus adopt the LTE rate to produce the main results in the present paper. In most cases, we find a constant “floor” of the H$_2$ fraction at $2 \times 10^{-7}$ owing to the H$_2$ formation via the H$_2^+$ channel (Fig. 4b).